

Enhancement of large magnetoresistances in ruthenocuprates by Ta substitution†

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An unexpected enhancement of the large negative magnetoresistance (MR) observed in $\text{RuSr}_2\text{Nd}_{0.95}\text{Y}_{0.15}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$ up to -47% at 4 K and 9 T is evidenced upon dilution of the Ru magnetic order by substitution of Ta for Ru; this enhancement of $-\text{MR}$ scales with the cell volume.

The phenomena of high-temperature superconductivity (HTSC) and colossal magnetoresistance (CMR) have been much studied in recent years. Both phenomena are observed in doped transition metal oxides; HTSC is observed in layered cuprates such as $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ ¹ between a Cu oxidation state of 2.06–2.30; and CMR is established in perovskite manganites such as $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ for Mn oxidation states of 3.2–3.5.²

Ruthenocuprates, containing distinct RuO_2 and CuO_2 layers, are an important class of multifunctional electronic materials that can display both properties. Coexisting weak ferromagnetism and superconductivity has previously been established in several 1212 ($\text{RuSr}_2\text{RCu}_2\text{O}_8$)^{3,4} and 1222 ($\text{RuSr}_2(\text{R,Ce})_2\text{Cu}_2\text{O}_{10-\delta}$) type materials,^{5–7} where R = Sm, Eu or Gd. We have recently reported that a new 1222 series, $\text{RuSr}_2\text{Nd}_{1.8-x}\text{Y}_{0.2}\text{Ce}_x\text{Cu}_2\text{O}_{10-\delta}$, exhibits large negative magnetoresistances (change of electrical resistivity ρ in an applied magnetic field H , defined as $\text{MR} = (\rho(H) - \rho(0))/\rho(0)$) for hole dopings $0.01 < p < 0.06$ (the equivalent Cu oxidation states are $2 + p$).^{8,9} Weak ferromagnetism arises due to Ru spin ordering below $T_{\text{Ru}} \sim 70\text{--}180$ K which induces long-range Cu spin ordering at $T_{\text{Cu}} \sim 20\text{--}110$ K. MR diverges to large negative values below T_{Cu} , for R = Nd, Y as the application of a magnetic field results in a canting of the Ru and Cu spins into a ferromagnetic arrangement in the ab plane. In this paper we will demonstrate an unexpected enhancement of the $-\text{MR}$ upon dilution of the magnetic order in the RuO_2 layer by substitution of Ta^{5+} for Ru^{5+} .

The 1222 structure consists of metal oxide layers in the repeat sequence $\text{RuO}_2.\text{SrO}.\text{CuO}_2.(R,\text{Ce}).\text{O}_{2-\delta}.(R,\text{Ce}).\text{CuO}_2.\text{SrO}$. A series of polycrystalline ceramic samples, $\text{Ru}_{1-x}\text{Ta}_x\text{Sr}_2\text{Nd}_{0.95}\text{Y}_{0.15}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$ ($0 < x < 0.4$) were prepared to investigate the effect of diluting the Ru magnetic order on the magnetotransport. Pelleted

mixtures of Nd_2O_3 , Y_2O_3 , RuO_2 , Ta_2O_5 , CuO , CeO_2 and SrCO_3 powders were repeatedly sintered at 1025°C and furnace cooled in air at a rate of 2°C min^{-1} . It is difficult to prepare single phase Nd-based 1222 materials as these only form over a $\sim 10^\circ\text{C}$ synthesis window around 1025°C , and a small degree of Y substitution was needed to produce samples with a high degree ($>95\%$) of phase purity. X-Ray diffraction patterns demonstrated that tetragonal $I4/mmm$ 1222 phases were formed in all cases (Fig. 1). An increase in both a and c is observed, from 3.8546(2) and 28.5610(9) Å to 3.8644(3) and 28.6405(9) Å, respectively, upon increasing x from 0 to 0.4, in accordance with the substitution of the larger Ta^{5+} ion for Ru^{5+} .

Magnetisations were recorded on a Quantum Design SQUID magnetometer in a 100 Oe magnetic field. As in previous samples,^{8,9} a weak ferromagnetic transition of the Ru spins is observed, followed by low temperature antiferromagnetic order of the Cu spins. The Ru spin transition, T_{Ru} , was determined by extrapolating the maximum ($-dM/dT$) slope to zero magnetization. Likewise the Cu antiferromagnetic transition was established from the temperature of the maximum magnetisation recorded in a field of 100 Oe after zero-field cooling.^{8,9} A reduction of both T_{Ru} and T_{Cu} , from 144 and 73 K to 102 and 55 K, respectively, is evidenced with increasing x , in keeping with the dilution of magnetic $4d^3 \text{Ru}^{5+}$ by diamagnetic Ta^{5+} . Similar results are observed upon substitution of Sn^{4+} and Nb^{5+} for Ru in $\text{RuSr}_2\text{GdCu}_2\text{O}_8$ ^{10,11} and $\text{RuSr}_2\text{Eu}_{1.2}\text{Ce}_{0.8}\text{Cu}_2\text{O}_{10-\delta}$.¹² The Ru spin order induces long-range antiferromagnetism in the CuO_2 plane⁹ (this long-range antiferromagnetic order is not normally observed for $p > 0.02$ in HTSC), so the reduction of T_{Ru} by Ta^{5+} substitution results in a decrease in T_{Cu} .

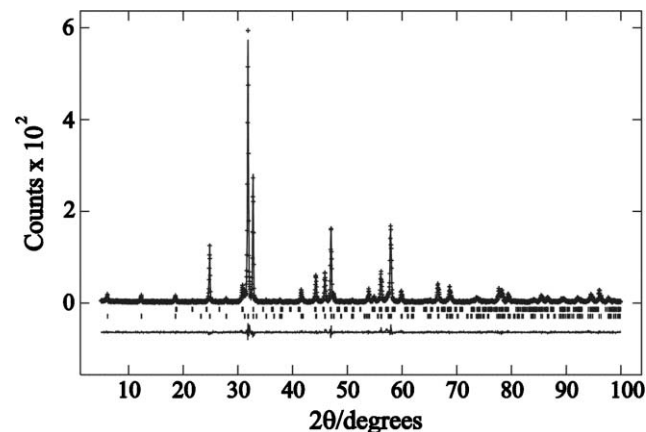


Fig. 1 Rietveld refinement fit to the powder X-ray diffraction pattern of $\text{Ru}_{0.6}\text{Ta}_{0.4}\text{Sr}_2\text{Cu}_2\text{Nd}_{0.95}\text{Y}_{0.15}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$; lower and upper reflection markers correspond to $\text{Ru}_{0.6}\text{Ta}_{0.4}\text{Sr}_2\text{Cu}_2\text{Nd}_{0.95}\text{Y}_{0.15}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$ and a trace of $\text{Sr}_2\text{Nd}(\text{Ta},\text{Ru})\text{O}_6$, respectively.

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† Electronic supplementary information (ESI) available: Table S1: Refined cell parameters and magnetic transition temperatures for the $\text{Ru}_{1-x}\text{Ta}_x\text{Sr}_2\text{Nd}_{0.95}\text{Y}_{0.15}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$ solid solutions. Fig. S1: Temperature variation of zero-field cooled magnetic susceptibility for the $\text{Ru}_{1-x}\text{Ta}_x\text{Sr}_2\text{Nd}_{0.95}\text{Y}_{0.15}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$ solid solutions. See DOI: 10.1039/b617872c

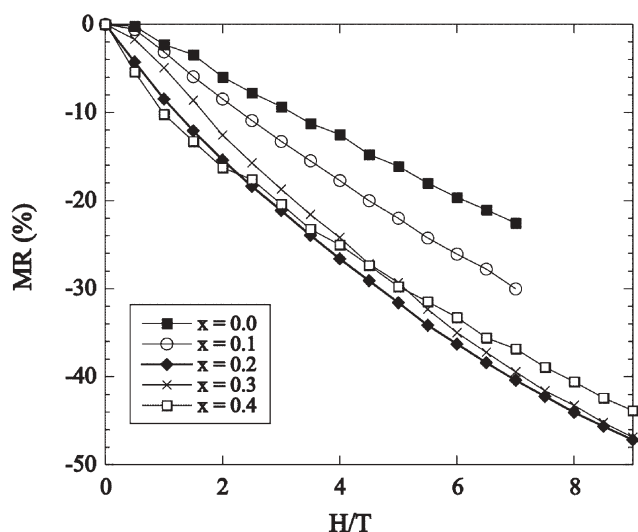


Fig. 2 The field evolution of MR for $\text{Ru}_{1-x}\text{Ta}_x\text{Sr}_2\text{Nd}_{0.95}\text{Y}_{0.15}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$ recorded at 4 K.

Resistances were measured on sintered polycrystalline bars on a Quantum Design physical property measurement system up to $H = 7$ or 9 T (Fig. 2). All samples are semiconducting. An unexpected enhancement of $-\text{MR}$ is observed as the Ta concentration increases; $\text{MR}_{7\text{T}}(4\text{ K})$ rises from -23 to -40% as x increases from 0 to 0.2. $-\text{MR}$ does not saturate in fields up to 9 T and the value of $\text{MR}_{9\text{T}}(4\text{ K}) = -47\%$ for $x = 0.2$ is the largest reported value for ruthenocuprates so far. A slight reduction in the high field $-\text{MR}$ is observed for higher Ta substitutions.

The large magnetoresistances of 1222 ruthenocuprates arise from hole transport in the antiferromagnetically ordered CuO_2 planes. As the Cu spin order is induced by Ru spin order, it is surprising that dilution of the Ru spin order with concomitant reductions of T_{Ru} and T_{Cu} results in a large enhancement of the magnetotransport properties up to 20% Ta substitution.

The likely explanation for the enhanced $-\text{MR}$ in $\text{Ru}_{1-x}\text{Ta}_x\text{Sr}_2\text{Nd}_{0.95}\text{Y}_{0.15}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$ lies in the sensitivity of Ru-1222 materials to lattice effects. Previous results⁹ have shown an approximately linear change from positive to negative MR as the R^{3+} size increases in the series $\text{RuSr}_2\text{R}_{1,1}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$. From analysis of the temperature dependence of MR, the negative regime is assigned to two-dimensional transport within the CuO_2 planes, whereas hole conduction between planes gives rise to positive MR. Fig. 3 shows that the MR's for $\text{Ru}_{1-x}\text{Ta}_x\text{Sr}_2\text{Nd}_{0.95}\text{Y}_{0.15}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$ up to $x = 0.2$ follow the same linear trend as the former $\text{RuSr}_2\text{R}_{1,1}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$ series with unit cell volume (equivalent correlations are observed with the a or c cell parameters, which vary monotonically across the two series). This suggests that the magnetotransport becomes more 'two-dimensional' and negative as the cell expands, due to some combination of the increasing separation of the metal oxide layers, and the flattening of the CuO_2 planes. Thus, the size effect of substituting Ta^{5+} for Ru^{5+} outweighs the electronic effects up to $x = 0.2$. Above this, $-\text{MR}$ decreases and no longer follows the linear trend, showing that the dilution effects on the magnetic order become significant, so that $x_c = 0.2$ is the critical dilution limit for Ru magnetism. Neutron diffraction will be needed to determine whether long-range order of Ru and/or Cu spins

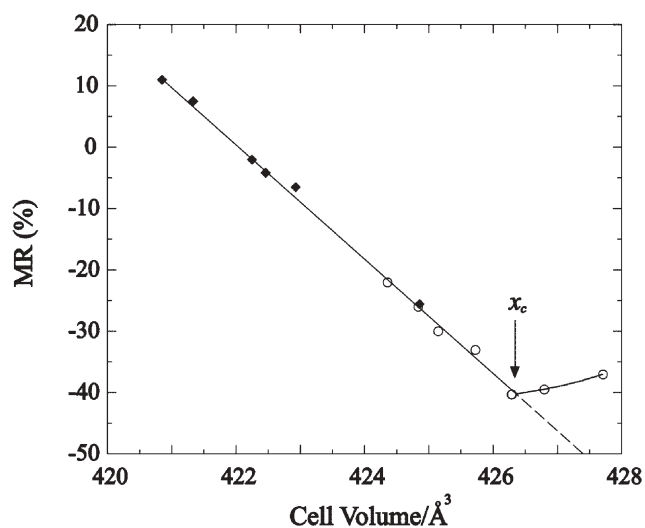


Fig. 3 Variation of $\text{MR}_{7\text{T}}(4\text{ K})$ with cell volume for $\text{RuSr}_2\text{R}_{1,1}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$ ($\text{R} = \text{Gd}, \text{Eu}, \text{Eu}_{0.9}\text{Y}_{0.2}, \text{Sm}, \text{Sm}_{0.9}\text{Y}_{0.2}, \text{Nd}_{0.9}\text{Y}_{0.2}$; diamonds) and $\text{Ru}_{1-x}\text{Ta}_x\text{Sr}_2\text{Nd}_{0.95}\text{Y}_{0.15}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$ ($x = 0, 0.05, 0.10, 0.15, 0.20, 0.30, 0.40$; circles). A deviation from linearity is observed above $x_c = 0.2$.

changes at x_c , and study of higher x materials will be interesting to follow the evolution of magnetotransport properties.

In conclusion, we find that the large $-\text{MR}$ previously observed in underdoped 1222 type ruthenocuprates can be enhanced by Ta substitution in the new series $\text{Ru}_{1-x}\text{Ta}_x\text{Sr}_2\text{Nd}_{0.95}\text{Y}_{0.15}\text{Ce}_{0.9}\text{Cu}_2\text{O}_{10-\delta}$, as the size effect outweighs the dilution of the magnetic RuO_2 layers up to a critical concentration of $x_c = 0.2$. Substitution of even larger cations into underdoped 1222 ruthenocuprates could result in further enhancements of the low temperature magnetoresistance of these complex, multifunctional electronic materials.

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